## Crystal Structure of Dilead Tritellurate(Iv)

By John C. Dewan and Anthony J. Edwards,* Chemistry Department, University of Birmingham, P.O. Box 363, Birmingham B15 2TT<br>Gordon R. Jones and Ian M. Young, Royal Signals and Radar Establishment, Great Malvern, Worcestershire

The crystal structure of the title compound has been determined by the heavy-atom method from 1059 reflections, measured with a diffractorneter and refined by full-matrix least-squares methods to $R 0.105$. Crystals are orthorhombic, space group Amam, with unit-cell dimensions $a=18.83(2), b=7.14(1), c=19.37(2) \AA$, and $Z=12$. In the structure two of the crystallographically independent tellurium atoms are linked into endless chains parallel to the $a$ axis by bridging oxygen atoms, and the other two are linked into isolated $\left[\mathrm{Te}_{3} \mathrm{O}_{8}\right]^{4-}$ ions. The co-ordination of the tellurium atoms can be considered to be based on a tetrahedral arrangement of three oxygen atoms and a lone pair of electrons, or a trigonal-bipyramidal arrangement of four oxygen atoms and the lone pair. The irregular co-ordination of the lead atoms can also be attributed to lone-pair effects.

Dilead tritellurate(iv) was first grown in large singlecrystal form by Robertson et al. ${ }^{1}$ who supplied the material for this study. The pale yellow transparent crystal shows a tendency to cleave on a major growth face, and was shown to be biaxial with a 2 V of $c a .72^{\circ}$. Only a few compounds of this stoicheiometry have been reported in the literature, and since the lead compound did not appear to be related to any previously observed materials the crystal-structure analysis was undertaken.

## EXPERIMENTAL

Approximate cell dimensions were determined from oscillation, Weissenberg, and precession photographs. The
photographs indicated orthorhombic symmetry and systematic absences corresponded to the space groups Amam , $A 2_{1} a m$, or $A m a 2$.

Final cell dimensions were measured on a Stoe two-circle computer-controlled diffractometer with graphite-monochromated Mo- $K_{\alpha}$ radiation and a scintillation counter. Intensity data were collected from a crystal fragment of dimensions $0.12 \times 0.08 \times 0.30 \mathrm{~mm}$ mounted about the direction of elongation (c). The $\omega$-scan technique was employed with a stepping interval of $0.01^{\circ}$ and a step time of 0.25 s ; this corresponds to a scan rate of $2.4^{\circ} \mathrm{min}^{-1}$. For layers $h k 0-6, \Delta \omega$ was taken as $2.0^{\circ}$, and for the seventh and higher layers (equi-inclination angle, $\mu>7^{\circ}$ ) a variable scan range was employed, $\Delta \omega$ being calculated
${ }^{1}$ D. S. Robertson, N. Shaw, and I. M. Young, J. Phys. (D), 1976, 9, 1257.
from $\left[A+\left(B \sin \mu / \tan \phi^{\prime}\right)\right]^{\circ}$, where $2 \phi^{\prime}$ is the azimuth angle ${ }^{2}$ and $A$ and $B$ were assigned values of 1.0 and 0.5 , respectively. The background for all the reflections was measured for 10 s at each end of the scan. Of 1355 reflections scanned within the range $0.1<(\sin \theta) / \lambda<0.65$, 1059, for which $I>3 \sigma(I)$, were considered to be observed and were used in the structure solution and refinement. Data were corrected for Lorentz and polarisation factors as well as for absorption, using a numerical integration method.
Crystal Data.- $\mathrm{O}_{8} \mathrm{~Pb}_{2} \mathrm{Te}_{3}, \quad M=925.2$, Orthorhombic, $a=18.83(2), b=7.14(1), c=19.37(2) \AA, U=2604.2 \AA^{3}$, $D_{\mathrm{m}}=7.16 \times 10^{3} \mathrm{~kg} \mathrm{~m}^{-3}, Z=12, D_{\mathrm{c}}=7.08, F(000)=$ 4 608. Systematic absences: $h k l$ when $k+l$ is odd, $h 0 l$ when $h$ is odd, space group $\operatorname{Amam}$ ( $D_{2 h}^{17}$, in a non-standard setting), $A 2_{1} a m$ ( $C_{2 v}^{12}$, in a non-standard setting), or $A m a 2$ $\left(C_{2 v}^{16}\right)$; Amam established as a result of the structure analysis. Mo- $K_{\alpha}$ radiation ( $\lambda 0.7107 \AA, \mu=465.6 \mathrm{~cm}^{-1}$ ).

Structure Determination.-The Patterson map was successfully solved in the centrosymmetric space group Amam yielding the co-ordinates of the two independent Pb atoms. Subsequent difference-Fourier maps were used to locate the Te and O atoms. The choice of space group has been vindicated by the successful solution and refinement of the structure. In addition, electrical measurements indicated the presence of a centre of symmetry. ${ }^{2}$ Initial refinement was by full-matrix least-squares methods with layer scale factors refined separately, and all the atoms vibrating isotropically. Refinement was continued with layer scale factors held constant and with the introduction of anisotropic thermal parameters, for all the Pb and Te atoms, of the form $\exp \left[-2 \pi^{2}\left(U_{11} h^{2} a^{* 2}+\ldots+2 U_{12} h k a^{*} b^{*}+\ldots\right)\right]$. Scattering factors employed were those for the neutral atoms, ${ }^{3}$ with corrections for the effects of anomalous dispersion ${ }^{4}$ applied to those for Pb and Te . The weighting scheme used in the final cycles was ${ }^{5} w=\left[\sigma^{2}\left(F_{0}\right)+\right.$ $\left.\left(0.07\left|F_{0}\right|\right)^{2}\right]^{-1}$, chosen to give approximately constant values of $\Sigma w\left(\left|F_{\mathrm{o}}\right|-\left|F_{\mathrm{c}}\right|\right)^{2}$ in groups of increasing $\left|F_{0}\right|$ and $(\sin \theta /) \lambda$. The refinement of the isotropic secondary extinction co-

Table 1
Final atomic positional parameters with estimated standard deviations in parentheses

| Atom | $x / a$ | $y / b$ | $z / c$ |
| :---: | :---: | :---: | :---: |
| $\mathrm{Pb}(1)$ | $0.4098(1)$ | 0.368 6(3) | 0.307 4(1) |
| $\mathrm{Pb}(2)$ | $\frac{1}{4}$ | $0.5101(5)$ | $0.1916(2)$ |
| $\mathrm{Te}(1)$ | $\frac{1}{4}$ | $0.1107(10)$ | $\frac{3}{2}$ |
| $\mathrm{Te}(2)$ | $0.0730(3)$ | 0.2979 (8) | $\frac{1}{2}$ |
| $\mathrm{Te}(3)$ | $0.4164(2)$ | 0.844 7(5) | $0.3865(2)$ |
| $\mathrm{Te}(4)$ | $\frac{1}{4}$ | 0.564 4(7) | $0.3879(3)$ |
| $\mathrm{O}(1)$ | $0.1314(26)$ | 0.0825 (74) | $\frac{1}{2}$ |
| $\mathrm{O}(2)$ | $\frac{1}{2}$ | 1 | $0.4167(31)$ |
| $\mathrm{O}(3)$ | 0.468 6(19) | $0.7139(56)$ | $0.3267(19)$ |
| $\mathrm{O}(4)$ | $0.3189(18)$ | $0.6807(50)$ | $0.3367(18)$ |
| $\mathrm{O}(5)$ | $0.3807(22)$ | $1.0419(60)$ | $0.3239(20)$ |
| $\mathrm{O}(6)$ | $0.1131(26)$ | 0.4323 (74) | $0.4309(23)$ |
| $\mathrm{O}(7)$ | $\frac{1}{4}$ | -0.0525(111) | $0.4236(36)$ |
| $\mathrm{O}(8)$ | $\frac{1}{4}$ | 0.3737 (103) | $0.3343(36)$ |

In the non-standard setting of the space group, general equivalent positions are: $x, y, z ; \bar{x}, \bar{y}, \bar{z} ; \bar{x}, \bar{y}, z ; x, y, \bar{z}$; $\frac{1}{2}-x, y, z ; \frac{1}{2}+x, \bar{y}, \bar{z} ; \frac{1}{2}-x, y, \bar{z} ;$ and $\frac{1}{2}+x, \bar{y}, z$.
efficient, ${ }^{6} g$, converged at $5.7(4)$. No parameter shift, at convergence, was greater than $0.001 \sigma$ and a final diffrence-

[^0]${ }^{3}$ D. T. Cromer and J. B. Mann, Acta Cryst., 1968, A24, 321.

Fourier map showed no significant features. The final $R$ for all the 1059 observed amplitudes is $10.5 \%$. The high $R$ value for this structure is probably a result of its high lead content and the difficulty of defining the crystal for absorption corrections.

Table 2
Interatomic distances $(\AA)$ and angles ( ${ }^{\circ}$ ) with estimated standard deviations in parentheses
(a) Distances

| $\mathrm{Te}(1)-\mathrm{O}(1)$ | 2.24(5) | $\mathrm{Te}(2)-\mathrm{O}(1)$ | 1.89(5) |
| :---: | :---: | :---: | :---: |
| $\mathrm{Te}(1)-\mathrm{O}(7)$ | 1.88(7) | $\mathrm{Te}(2)-\mathrm{O}(6)$ | 1.81 (5) |
| $\mathrm{Te}(3)-\mathrm{O}(2)$ | 2.01(2) | $\mathrm{Te}(2) \cdots \mathrm{O}\left(2^{\mathrm{I}}\right)$ | 3.00 (3) |
| $\mathrm{Te}(3)-\mathrm{O}(4)$ | 2.38(3) | $\mathrm{Te}(4)-\mathrm{O}(4)$ | 1.83(3) |
| $\mathrm{Te}(3)-\mathrm{O}(3)$ | 1.78(4) | $\mathrm{Te}(4)-\mathrm{O}(8)$ | 1.71 (7) |
| $\mathrm{Te}(3)-\mathrm{O}(5)$ | 1.98(4) | $\mathrm{Te}(4) \cdots \mathrm{O}(6)$ | 2.87(5) |
| $\mathrm{Te}(3) \cdots \mathrm{O}\left(1^{\text {II }}\right.$ ) | 2.92(5) | $\mathrm{Te}(4) \cdots \mathrm{O}\left(\mathbf{7}^{\text {IV }}\right.$ ) | 2.82(7) |
| $\mathrm{Te}(3) \cdots \mathrm{O}\left(6{ }^{\text {III }}\right.$ ) | 3.12(6) | $\mathrm{Pb}(1) \cdots \mathrm{O}(3)$ | 2.73(4) |
| $\mathrm{Te}(3) \cdots \mathrm{O}\left(7^{\text {IV }}\right.$ ) | 3.30 (7) | $\mathrm{Pb}(1) \cdots \mathrm{O}(4)$ | 2.87(4) |
| $\mathrm{Pb}(2) \cdots \mathrm{O}(4)$ | 3.33(3) | $\mathrm{Pb}(1) \cdots \mathrm{O}(8)$ | $3.05(7)$ |
| $\mathrm{Pb}(2) \cdots \mathrm{O}(8)$ | 2.93(7) | $\mathrm{Pb}(1) \cdots \mathrm{O} 3^{\text {V }}$ | 2.39(4) |
| $\mathrm{Pb}(2) \cdots \mathrm{O}\left(4^{\text {VI }}\right)$ | 2.74(4) | $\mathrm{Pb}(1) \cdots \mathrm{O}\left(3^{\mathrm{VI}}\right)$ | 3.03(4) |
| $\mathrm{Pb}(2) \cdots \mathrm{O}\left(5^{\text {VI }}\right)$ | 2.49(4) | $\mathrm{Pb}(1) \cdots \mathrm{O}\left(5^{\mathrm{VII}}\right)$ | 2.42 (4) |
| $\mathrm{Pb}(2) \cdots \mathrm{O}\left(7^{\text {VIII }}\right)$ | 2.28(7) | $\mathrm{Pb}(1) \cdots \mathrm{O}\left(5^{\mathrm{VI}}\right)$ | 2.88(4) |
| $\mathrm{Pb}(2) \cdots \mathrm{O}(8 \mathrm{VIII})$ | 2.64(7) | $\mathrm{Pb}(1) \cdots \mathrm{O}\left(6{ }^{\text {III }}\right.$ ) | 2.47 (4) |
| $\mathrm{O}(1) \cdots \mathrm{O}(7)$ | 2.85(6) | $\mathrm{O}(1) \cdots \mathrm{O}(6)$ | 2.85 (7) |
| $\mathrm{O}(7) \cdots \mathrm{O}\left(7^{\text {IX }}\right.$ ) | 2.96(10) | $\mathrm{O}(6) \cdots \mathrm{O}\left(6^{1 \mathrm{X}}\right)$ | 2.68 (6) |
| $\mathrm{O}(3) \cdots \mathrm{O}(5)$ | 2.87(6) | $\mathrm{O}(2) \cdots \mathrm{O}(3)$ | 2.75 (5) |
| $\mathrm{O}(4) \cdots \mathrm{O}(5)$ | 2.84(6) | $\mathrm{O}(2) \cdots \mathrm{O}(5)$ | 2.89 (6) |
| $\mathrm{O}(4) \cdots \mathrm{O}(3)$ | 2.84(5) | $\mathrm{O}(4) \cdots \mathrm{O}(8)$ | 2.55 (7) |

(b) Angles

| $\mathrm{O}(1)-\mathrm{Te}(1)-\mathrm{O}\left(1^{1 \mathrm{I}}\right)$ | 169.7(1.9) |
| :---: | :---: |
| $\mathrm{O}(7)-\mathrm{Te}(1)-\mathrm{O}\left(7^{\mathbf{1 2}}\right)$ | 103.6(3.3) |
| $\mathrm{Te}(1)-\mathrm{O}(1)-\mathrm{Te}(2)$ | 120.4(2.5) |
| $\mathrm{O}(1)-\mathrm{Te}(1)-\mathrm{O}(7)$ | 86.8(0.9) |
| $\mathrm{O}(4)-\mathrm{Te}(4)-\mathrm{O}(8)$ | 91.9(2.0) |
| $\mathrm{O}(4)-\mathrm{Te}(4)-\mathrm{O}\left(4{ }^{\text {III }}\right.$ ) | $90.2(1.5)$ |
| $\mathrm{Te}(3)-\mathrm{O}(4)-\mathrm{Te}(4)$ | 123.3(1.7) |
| $\mathrm{O}\left(3^{\mathrm{V}}\right)-\mathrm{Pb}(1)-\mathrm{O}\left(5^{\mathrm{vII}}\right)$ | 87.6(1.4) |
| $\mathrm{O}\left(3^{\mathrm{V}}\right)-\mathrm{Pb}(1)-\mathrm{O}\left(6^{\text {III }}\right)$ | 93.5 (1.4) |
| $\mathrm{O}\left(5^{\text {VII }}\right)-\mathrm{Pb}(1)-\mathrm{O}(6 \mathrm{III})$ | 90.5(1.5) |
| $\mathrm{O}\left(5^{\mathrm{VI}}\right)-\mathrm{Pb}(2)-\mathrm{O}\left(5^{\mathrm{xI}}\right)$ | 162.6(1.3) |
| $\mathrm{O}(1)-\mathrm{Te}(2)-\mathrm{O}(6)$ | 100.9(1.9) |
| $\mathrm{O}(6)-\mathrm{Te}(2)-\mathrm{O}\left(6^{\text {TX }}\right.$ ) | 95.2(2.2) |
| $\mathrm{O}(2)-\mathrm{Te}(3)-\mathrm{O}(3)$ | 92.6(1.6) |
| $\mathrm{O}(2)-\mathrm{Te}(3)-\mathrm{O}(4)$ | 172.6(1.8) |
| $\mathrm{O}(2)-\mathrm{Te}(3)-\mathrm{O}(5)$ | 93.0(1.6) |
| $\mathrm{O}(3)-\mathrm{Te}(3)-\mathrm{O}(5)$ | 99.3 (1.7) |
| $\mathrm{O}(4)-\mathrm{Te}(3)-\mathrm{O}(5)$ | 80.8(1.5) |
| $\mathrm{O}(4)-\mathrm{Te}(3)-\mathrm{O}(3)$ | 84.5(1.4) |
| $\mathrm{Te}(3)-\mathrm{O}(2)-\mathrm{Te}\left(3^{\mathrm{x}}\right)$ | 146.2(3.3) |
| $\mathrm{O}\left(7^{\mathrm{VIII}}\right)-\mathrm{Pb}(2)-\mathrm{O}\left(8^{\mathrm{VIII}}\right)$ | 90.4(2.5) |


| I $\frac{1}{2}-x,-1+y, z$ | II $\frac{1}{2}-x, 1+y, z$ |
| :--- | ---: |
| III $\frac{1}{2}-x, y, z$ | VI $x, 1+y, z$ |
| V $1-x, 1-y, z$ | VI $x,-\frac{1}{2}+y, \frac{1}{2}-z$ |
| VII $x,-1+y, z$ | VIII $x, \frac{1}{2}+y, \frac{1}{2}-z$ |
| IX $x, y, 1-z-\frac{3}{2}$ | X $1-x, 2-y, z$ |
| XI $\frac{1}{2}-x, \frac{3}{2}-y, \frac{1}{2}-z$ |  |

Most of the computations were carried out on the CDC 7600 at the University of Manchester Regional Computer Centre using the ' $X$-RAY ' 72 ' system of programs; SHELX-76 on the ICL 1906A at the University of Birmingham Computer Centre was used for the absorption corrections. Final atomic co-ordinates with estimated standard deviations are listed in Table 1, interatomic distances and angles in Table 2. Observed and calculated structure amplitudes ( $\times 10$ ) and anisotropic thermal parameters are listed in Supplementary Publication No. SUP 22306 ( 5 pp .).*
${ }^{4}$ D. T. Cromer, Acta Cryst., 1965, 18, 17.
${ }^{5}$ G. H. Stout and L. H. Jensen, ' $X$-Ray Structure Determination,' Macmillan, London, 1969, p. 456.
${ }^{6}$ A. C. Larson, Acta Cryst., 1967, 23, 664.

## RESULTS AND DISCUSSION

The structure is illustrated in Figure 1 and can best be described in terms of separate $\mathrm{Pb}^{2+}$ cations and complex tellurate(iv) anions. The anions are of two types, those defined by $\mathrm{Te}(1)$ and $\mathrm{Te}(2)$ and those defined by $\mathrm{Te}(3)$ and $\mathrm{Te}(4)$. $\mathrm{Te}(1)$ in a four-fold and $\mathrm{Te}(2)$ in an eightfold special position, with $O(1), O(6)$, and $O(7)$ in eight-, sixteen-, and eight-fold positions respectively, give four separate $\left[\mathrm{Te}_{3} \mathrm{O}_{8}\right]^{4-}$ anions in the unit cell. $\mathrm{Te}(3)$ in a sixteen-fold and $\mathrm{Te}(4)$ in an eight-fold position, with their associated oxygen atoms, form endless-chain anions parallel to the $a$ axis, with eight $\left[\mathrm{Te}_{3} \mathrm{O}_{8}\right]$ units in the unit cell.

In the separate anion the central tellurium atom $\mathrm{Te}(\mathbf{1})$ forms two bonds to terminal oxygen atoms, $\mathrm{O}(7)$, of


Figure 1 Projection of $\frac{1}{4}$ of the unit cell down [010] showing the atom numbering
$1.88 \AA$ and two to bridging oxygen atoms, $O(1)$, of $2.24 \AA$. The co-ordination geometry can best be described in terms of a trigonal-bipyramidal arrangement, on valence-shell electron-pair repulsion (VSEPR) theory, ${ }^{7}$ with the lone pair occupying the fifth coordination position. The $\mathrm{O}(7)-\mathrm{Te}(1)-\mathrm{O}\left(7^{\mathrm{I}}\right)$ angle of $103.6^{\circ}$ and the $\mathrm{O}(1)-\mathrm{Te}(1)-\mathrm{O}\left(1^{\mathrm{III}}\right)$ angle of $169.7^{\circ}$ are in line with the theoretical angles of 120 and $180^{\circ}$, when lone pair-bond pair interactions are taken into account.

The outer tellurium atoms $\mathrm{Te}(2)$ form two bonds to terminal oxygen atoms $\mathrm{O}(6)$ of $1.81 \AA$ and one bond to the bridging oxygen atom $O(1)$ of $1.89 \AA$. The pyramid arrangement of the three oxygen atoms, with an average $\mathrm{O}-\mathrm{Te}-\mathrm{O}$ angle of $99^{\circ}$, can be correlated with a tetrahedral arrangement around tellurium, including the lone pair of electrons.
The $\mathrm{Te}(1)-\mathrm{O}(1)-\mathrm{Te}(2)$ bridge is asymmetric, and $\mathrm{O}(1)$ is more closely associated with $\mathrm{Te}(2)$, which suggests a contribution to the structure from the form $\mathrm{TeO}_{3}{ }^{2--}$

[^1]$\mathrm{TeO}_{2}-\mathrm{TeO}_{3}{ }^{2-}$ with two pyramidal ions linked by a neutral $\mathrm{TeO}_{2}$ molecule.
In the chain anion $\mathrm{Te}(4)$ has three short bonds to oxygen of $1.83 \AA(\times 2)$ to $O(4)$ and $1.71 \AA$ to $O(8)$. The distance of $1.71 \AA$ is extremely short, but the estimated standard deviation is high, and it is not significantly different from the other $\mathrm{Te}-\mathrm{O}$ terminal bonds. The $\mathrm{O}-\mathrm{Te}-\mathrm{O}$ angles average $91^{\circ}$, giving a similar pyramidal co-ordination to that of $\mathrm{Te}(2)$, again on VSEPR theory based on a tetrahedral arrangement including the lone pair. $\mathrm{Te}(3)$ has two short terminal bonds to $O(3)$ and $O(5)$ of 1.78 and $1.98 \AA$, and a short symmetric bridge bond to $O(2)$ of $2.01 \AA$, with an average $\mathrm{O}-\mathrm{Te}-\mathrm{O}$ angle of $95^{\circ}$, to give a similar pyramidal co-ordination to those of $\mathrm{Te}(4)$ and $\mathrm{Te}(\mathbf{2})$.

The chain is completed by the long asymmetric bridge bond $\mathrm{Te}(3)-\mathrm{O}(4)$ of $2.38 \AA$. This gives a four-co-ordinate arrangement around $\mathrm{Te}(3)$ which is similar to that around $\mathrm{Te}(\mathbf{1})$, and can again be described in terms of a trigonal-bipyramidal co-ordination, including the lone pair of electrons. The $\mathrm{O}(3)-\mathrm{Te}(3)-\mathrm{O}(5)$ and $\mathrm{O}(2)^{-}$ $\mathrm{Te}(3)-\mathrm{O}(4)$ angles of 99.5 and $173.6^{\circ}$ are very similar to the corresponding angles at $\mathrm{Te}(\mathbf{1})$ described above. The asymmetric bridge again suggests a contribution to the structure from the ionic form $\left[\mathrm{TeO}_{3}\right]^{2-}+\left[\mathrm{Te}_{2} \mathrm{O}_{5}\right]^{2-}$, with the ions linked through the weak bond.

The tellurate(iv) anions in the present structure thus contain tellurium atoms which can be described as threeand four-co-ordinated by oxygen atoms. These two configurations have been found previously; for example, teinite, ${ }^{8} \mathrm{CuTeO}_{3} \cdot 2 \mathrm{H}_{2} \mathrm{O}$, contains a three-co-ordinate pyramidal unit with an average $\mathrm{Te}-\mathrm{O}$ distance of $1.95 \AA$ and average angles of $95^{\circ}$. Both the paratellurite ${ }^{9}$ and tellurite ${ }^{10}$ forms of $\mathrm{TeO}_{2}$ contain four-co-ordinate units, with very similar geometries, based on the trigonalbipyramidal model, to those described here. Average values for $\mathrm{Te}-\mathrm{O}$ (axial) are 2.10 and for $\mathrm{Te}-\mathrm{O}$ (equatorial) $1.85 \AA$, and corresponding angles $c a .100$ and $150-180^{\circ}$.

Table 3
Parameters of some trigonal-bipyramidal co-ordinated $\mathrm{TeO}_{4}$ units

|  | $\mathrm{Te}-\mathrm{O}$ bond $/ \AA$ |  | $\begin{gathered} \mathrm{O}-\mathrm{Te}-\mathrm{O} \\ \text { Angle } /^{\circ} \end{gathered}$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Compound | apical | axial | apical | axial | Ref. |
| Theoretical model | 2.23 | 1.88 | 162 | 102 | 11 |
| $\mathrm{Zn}_{2} \mathrm{Te}_{3} \mathrm{O}_{8}[\mathrm{Te}(\mathrm{l})]$ | 2.10 | 1.83 | 175 | 106 | 13 |
| $\mathrm{TiTe}_{3} \mathrm{O}_{8}$ | 2.12 | 1.85 | 159 | 102 | $a$ |
| $\mathrm{Fe}_{2} \mathrm{Te}_{4} \mathrm{O}_{11}[\mathrm{Te}(4)]$ | 2.17 | 1.96 | 150 | 99 | 14 |
| $\mathrm{Pb}_{2} \mathrm{Te}_{3} \mathrm{O}_{8}[\mathrm{Te}(\mathrm{l})]$ | 2.24 | 1.88 | 170 | 104 | $b$ |
| [ $\mathrm{Te}(3)$ ] | 2.20 | 1.88 | 173 | 99 | $b$ |

${ }^{a}$ G. Meunier and J. Galy, Acta Cryst., 1971, B27, 602. ${ }^{b}$ This work.

Some values for both three- and four-co-ordinate species are given in Table 3, together with values from a theoretical model due to Galy et al. ${ }^{11}$

[^2]The asymmetric arrangement found for $\mathrm{Te}(\mathbf{3})$ has been described ${ }^{11}$ as $(3+1)$ co-ordination, and is a feature of a number of tellurate(rv) structures, e.g. denningite ${ }^{12}$


Figure 2 Projection down [100] of the co-ordination sphere of $\mathrm{Pb}(1)$ and down [010] of the co-ordination sphere of $\mathrm{Pb}(2)$, showing the atom numbering
$(\mathrm{Mn}, \mathrm{Ca}, \mathrm{Zn}) \mathrm{Te}_{2} \mathrm{O}_{5}$ with $\mathrm{Te}-\mathrm{O}$ distances of 2.04 and 2.36 , $\mathrm{Zn}_{2} \mathrm{Te}_{3} \mathrm{O}_{8}$ with distances ${ }^{13}$ of 1.98 and 2.41 , and $\mathrm{Fe}_{2} \mathrm{Te}_{4} \mathrm{O}_{11}$ with distances ${ }^{14}$ of 2.08 and $2.25 \AA$.
In addition to the near-neighbour co-ordination described above, there are some longer contacts to the tellurium atoms, a feature of structures containing $\mathrm{Te}^{\mathrm{tV}}$ bonded to either oxygen or fluorine. These long contacts have been included in a consideration of the total co-ordination by Brown, ${ }^{15}$ who has calculated bond valences for many examples and has tabulated structures in terms of three different configurations. These are based on a six-co-ordinate arrangement, with either two strong, two intermediate, and two weak bonds ( $\mathscr{A}$ ), or three strong and three weak bonds ( $\mathscr{C}$ ), or the transitional case where the two intermediate bonds are unequal ( $\mathscr{B}$ ).

In the present case the arrangements are not exactly comparable with Brown's configurations except for $\mathrm{Te}(4)$. Thus $\mathrm{Te}(1)$ has no further bonding [the nearest contacts to $O(6)$ and $O(8)$ are at $3.72 \AA$ ] and must be

[^3]considered four-co-ordinate (although it would be assigned to configuration $\mathscr{A})$. $\mathrm{Te}(2)$ has only two contacts to $\mathrm{O}(2)$ of $3.00 \AA$, and can be regarded as five-co-ordinate, best assigned to configuration $\mathscr{C}$ : $\mathrm{Te}(3)$ has contacts of $2.92 \AA$ to $\mathrm{O}(\mathrm{I}), 3.12 \AA$ to $\mathrm{O}(6)$, and $3.30 \AA$ to $O(7)$, and can therefore be considered to be seven-coordinate, although the interactions with $O(6)$ and $O(7)$ are obviously very weak. The arrangement is best assigned to configuration $\mathscr{A}$. $\mathrm{Te}(4)$ has contacts of $2.82 \AA$ to $\mathrm{O}(7)$ and $2.87 \AA(\times 2)$ to $\mathrm{O}(6)$ and is therefore six-co-ordinate and equivalent to configuration $\mathscr{C}$.
The lead atoms, although formally regarded as $\mathrm{Pb}^{2+}$ ions, also have a significant co-ordination by oxygen atoms. Thus $\mathrm{Pb}(\mathrm{l})$ has three nearest neighbours at an average distance of $2.42 \AA$, with an average $\mathrm{O}-\mathrm{Pb}-\mathrm{O}$ angle of $90.5^{\circ}$. This pyramidal arrangement can be considered on VSEPR theory in terms of a tetrahedral arrangement including the lone pair of electrons, similar to that for $\mathrm{Te}(2)$ and $\mathrm{Te}(4)$ above. A similar arrangement has been found in lead(II) oxide, ${ }^{16}$ lead(II) zinc silicate, ${ }^{17}$ and lead(II) germanate. ${ }^{18}$ There are five further contacts to the lead atom, as shown in Figure 2, giving a rather irregular arrangement, but grouped around the assumed lone-pair position.
For $\mathrm{Pb}(2)$ the situation is more difficult. However, the nearest neighbours are all on the same side of the lead atom, leaving space for the lone pair. Four oxygen atoms at an average distance of $2.48 \AA$, with the lone pair, form a distorted trigonal-bipyramidal arrangement with $\mathrm{O}-\mathrm{Pb}-\mathrm{O}$ (equatorial) of 90.4 and $\mathrm{O}-\mathrm{Pb}-\mathrm{O}$ (axial) of $162.6^{\circ}$. The $\mathrm{Pb}(2)-\mathrm{O}\left(7^{\mathrm{V}}\right)$ distance $(2.28 \AA)$ is extremely short but the estimated standard deviation is ligh and this is not considered significantly different


Figure 3 Projection of two adjacent cells along [001], showing the primary lead co-ordination polyhedra, in relation to the $\left[\mathrm{Te}_{3} \mathrm{O}_{8}\right]^{4-}$ chain network for atoms with $z$ co-ordinates between 0.065 and 0.195 . Key as in Figure 1
from the other nearest-neighbour distances. This type of co-ordination has also been found in the structures of
${ }^{17}$ C. T. Prewitt, E. Kirchner, and A. Preisinger, Z. Krist., 1967, 124, 115.
${ }_{18}$ M. I. Kay, R. E. Newnham, and R. W. Wolfe, Ferroelectrics, 1975, 9, 1 .
orthorhombic lead(II) oxide ${ }^{19}$ and lead(II) zinc silicate. ${ }^{17}$ The longer contacts to the lead atom again form a rather irregular arrangement and are grouped around the lone-pair position.
The nearest-neighbour polyhedra around $\mathrm{Pb}(1)$ and $\mathrm{Pb}(2)$ link to give $\mathrm{Pb}_{3} \mathrm{O}_{8}$ units which are illustrated in Figure 3. It can be seen from Figure 1 that the lead atoms are arranged with their lone pairs directed parallel to the $c$ axis, and there is clearly a weak structural ${ }^{10}$ M. I. Kay, Acta Cryst., 1961, 14, 80.
interaction in this direction, reflected in the long $\mathrm{Pb} \cdots \mathrm{O}$ contacts. This feature can be correlated with the pronounced tendency of the crystals to cleave on the $\{001\}$ face.

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